

## Entanglement-assisted Classical Capacity (Ch. 20)

Communication Resources: 1) Memoryless Channel  $(\mathcal{N}^{A \rightarrow B})^{\otimes n}$

2) Entanglement: Arbitrary state  $\phi^{A'B'}$

Encoder:  $\mathcal{E}_m^{A' \rightarrow A^n}(\phi^{A'B'})$

Decoder:  $\mathcal{L}_m^{B'B^n}$

Prob. error:  $p_e(m) = 1 - \text{Tr}(\mathcal{L}_m^{B'B^n} \mathcal{N}^{A' \rightarrow B^n}(\mathcal{E}_m^{A' \rightarrow A^n}(\phi^{A'B'})))$

Theorem (Bennett-Shor-Smolin-Thapliyal):

$$\sup \{R : \text{achievable}\} = I(\mathcal{N})$$

First observation:

All entanglement is equal (iid pure states)

Let  $\rho^{AB}$  be a pure state.

$$\rho^{AB} = H(A)_\rho [qq]$$

Entanglement concentration gives  $\rho^{AB} \geq H(A)_\rho [qq]$

The other direction works too.

Proof:

Achievability:

Main idea: Use the entanglement to construct a new channel.  
Use previous theorem (lower bound) about classical communication.

Review of lower bound:

$$\begin{aligned} \chi(\mathcal{N}) &= \sup_{\rho_{\bar{X}A}} \mathcal{I}(\bar{X}; \bar{B}) = \sup_{\text{cq state}} H(\bar{B}) - H(\bar{B}|\bar{X}) \\ &= \sup_{\{(p(x), \rho_x^{\bar{B}})\}} H(\bar{B})_{E_x \rho_x^{\bar{B}}} - E_x H(\bar{B})_{\rho_x^{\bar{B}}} \\ &\quad \uparrow \\ &\quad \mathcal{N}(\rho_x^A) \end{aligned}$$

Bottom line:

If an ensemble of channel outputs  $\bar{B}$  can be constructed then

$$H(\bar{B})_{E_x \rho_x} - E_x H(\bar{B})_{\rho_x} \text{ is achievable.}$$

(i.e. a code can be constructed randomly from the ensemble)

A special case:

$$\text{Suppose } \mathcal{I}(\mathcal{N}) = \mathcal{I}(X; B)_{\mathcal{N}^{A \rightarrow B}(\Phi^{XA})}$$

↑  
Bell state

(Recall,  $\mathcal{I}(\mathcal{N})$  is always optimized by a pure state.)  
 $\Rightarrow$  Pick input density  $\rho^A$  and consider its purification.  
 Assumption about is that  $\rho^A = \pi$  is optimal.

Let the entanglement resource be a Bell state  $|\Phi\rangle^{A'B'}$  with the same dimension as  $A$ .

Consider the following set of operations at the transmitter:

$$X^{A'}(x) Z^{A'}(z) |\Phi\rangle^{A'B'} \text{ then isomorphically map } A' \rightarrow A$$

Call the resulting state  $|\Phi_{xz}\rangle^{AB'}$  (this is the Bell basis)

Choose  $X$  and  $Z$  uniformly at random to produce an ensemble,  $\bar{B} = (B, B')$

$$1.) E_{x,z} \Phi_{x,z}^{AB'} = \pi^A \otimes \pi^{B'}$$

$$\Rightarrow E_{x,z} \mathcal{N}^{A \rightarrow B}(\Phi_{x,z}^{AB'}) = \mathcal{N}^{A \rightarrow B}(E \Phi_{x,z}^{AB'}) = \mathcal{N}^{A \rightarrow B}(\pi^A) \otimes \pi^{B'}$$

$$\Rightarrow H(\bar{B})_{E_{\rho_x}} = H(B, B')_{E_{x,z} \mathcal{N}^{A \rightarrow B}(\Phi_{x,z}^{AB'})} = H(B)_{\mathcal{N}(\pi)} + H(B')_{\pi}$$

$$2.) E_x H(\bar{B})_{\rho_x} = E_{x,z} H(B, B')_{\mathcal{N}^{A \rightarrow B}(\Phi_{x,z}^{AB'})}$$

Transpose result: Exercise 3.6.12

$$(M^A \otimes I^B) |\Phi\rangle^{AB} = (I^A \otimes (M^T)^B) |\Phi\rangle^{AB}$$

$$\Rightarrow |\Phi_{x,z}\rangle^{AB'} = Z^{TB'}(z) X^{TB'}(x) |\Phi\rangle^{AB'}$$

$$\begin{aligned} \Rightarrow \mathcal{N}^{A \rightarrow B}(\Phi_{x,z}^{AB'}) &= \mathcal{N}^{A \rightarrow B}(Z^{TB'}(z) X^{TB'}(x) \Phi_{x,z}^{AB'} X^{TB'}(x) Z^{TB'}(z)) \\ &= Z^{TB'}(z) X^{TB'}(x) \mathcal{N}^{A \rightarrow B}(\Phi_{x,z}^{AB'}) X^{TB'}(x) Z^{TB'}(z) \end{aligned}$$

$$\Rightarrow E_{x,z} H(B, B')_{\mathcal{N}^{A \rightarrow B}(\Phi_{x,z}^{AB'})} = H(B, B')_{\mathcal{N}^{A \rightarrow B}(\Phi)}$$

↑  
constant  $\forall x, z$

$$\text{Conclusion: } H(B)_{\mathcal{N}(\pi)} + H(B')_{\pi} - H(B, B')_{\mathcal{N}^{A \rightarrow B}(\Phi)} = I(B, B')_{\mathcal{N}^{A \rightarrow B}(\Phi_{x,z}^{AB'})}$$

is achievable.

$\Rightarrow I(N)$  is achievable by assumption.

Superdense coding is a special case with noise-free channel.

Generalize: (Idea 1)

Map Bell state to a type matching the capacity achieving input density.

This is like a constant composition code.

Don't know an easy way to complete the proof.

Generalization: (Idea 2)

Ensemble over super-symbols.

Entanglement resource is iid capacity achieving pure state  $(\rho^{A'B'})^{\otimes n}$

Use Schmidt decomposition and method of types.

$$|\rho\rangle^{A^n B^n} = \sum_{\mathbf{t}} \sqrt{p(\mathbf{t})} |\Phi_{\mathbf{t}}\rangle^{A^n B^n}$$

Construct ensemble from the following set of operations at the transmitter:

For each type let  $V_{\mathbf{t}}(x_{\mathbf{t}}, z_{\mathbf{t}}) = X_{\mathbf{t}}(x_{\mathbf{t}}) Z_{\mathbf{t}}(z_{\mathbf{t}})$

Pauli operators within type subspace.

$$\text{Let } U = \sum_{\mathbf{t}} (-1)^{b_{\mathbf{t}}} V_{\mathbf{t}}(x_{\mathbf{t}}, z_{\mathbf{t}})$$

This is a unitary operator parameterized by  $((b_{\mathbf{t}}, x_{\mathbf{t}}, z_{\mathbf{t}}))_{\mathbf{t}}$

Choose the parameters uniformly at random, apply  $U$  at transmitter, and map isomorphically to  $A^n$  to construct ensemble  $\rho_{(b_{\mathbf{t}}, x_{\mathbf{t}}, z_{\mathbf{t}})}^{A^n B^n} = (N_{x, \mathbf{I}}^{\otimes n}) (I^{A \rightarrow A^n} U^{A^n A^n} \rho^{A^n B^n} U^{\dagger} I^{\dagger})$

Verify that  $U$  is unitary:

A unitary transformation is equivalent to  $\{|\psi_i\rangle\} \rightarrow \{|\varphi_i\rangle\}$ ,  
where  $\{|\psi_i\rangle\}$  and  $\{|\varphi_i\rangle\}$  are both orthonormal bases.

$$\Rightarrow \sum_i |\varphi_i\rangle \langle \psi_i|$$

Furthermore,  $V_{\mathbf{t}}(x_{\mathbf{t}}, z_{\mathbf{t}}) = \sum_j |\varphi_{\mathbf{t},j}\rangle \langle \psi_{\mathbf{t},j}|$  where  $\{|\varphi_{\mathbf{t},j}\rangle\}$  and  $\{|\psi_{\mathbf{t},j}\rangle\}$  are orth. bases for type subspace.

Since, type subspaces are orth. and span  $A^n$ ,  $U$  is unitary.

\* Notice that  $U$  does not imply a measurement of the type and a conditional operation.

Claim 1:  $U^{A^n} \rho^{A^n B^n} = U^{\dagger B^n} \rho^{A^n B^n}$

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$\Rightarrow H(B^n, B'^n)_{\varphi_{(b_+, x_+, z_+)}}$   $= H(B^n, B'^n)_{\varphi^{B^n B'^n}}$

Proof:

$$U^{A^n} \rho^{A^n B^n} = \left( \sum_{\tau} (-1)^{b_{\tau}} V(x_{\tau}, z_{\tau}) \right)^{A^n} \left( \sum_{\tau} \sqrt{p(\tau)} |\Phi_{\tau}^{\pm}\rangle^{A^n B'^n} \right)$$

$$= \sum_{\tau} (-1)^{b_{\tau}} \sqrt{p(\tau)} V(x_{\tau}, z_{\tau})^{A^n} |\Phi_{\tau}^{\pm}\rangle^{A^n B'^n}$$

$$= \sum_{\tau} (-1)^{b_{\tau}} \sqrt{p(\tau)} V^{\tau}(x_{\tau}, z_{\tau})^{B^n} |\Phi_{\tau}^{\pm}\rangle^{A^n B'^n}$$

Claim 2:  $E \varphi_{(b_+, x_+, z_+)}^{B^n B'^n} = \sum_{\tau} p(\tau) \mathcal{N}^{A^n \rightarrow B^n}(\Pi_{\tau}^{A^n}) \otimes \Pi_{\tau}^{B^n}$

Not entangled but correlated based on the type.

$\Rightarrow H(E \varphi^{B^n B'^n}) = H(B^n)_{\mathcal{N}(\rho^{\otimes n})} + H(B'^n | T)_{\rho^{\otimes n}}$

$$= H(B^n)_{\mathcal{N}(\rho^{\otimes n})} + H(B'^n)_{\rho^{\otimes n}} - I(B'^n; T)_{\rho^{\otimes n}}$$

$$\geq \downarrow + \downarrow - \dim(B') \log n$$

Conclusion:

$$H(B^n)_{\mathcal{N}(\rho^{\otimes n})} + H(B'^n)_{\rho^{\otimes n}} - H(B^n, B'^n)_{\varphi^{\otimes n}} - \dim(B') \log n$$

$$= n \left( H(B)_{\mathcal{N}(\rho)} + H(B')_{\rho} - H(B, B')_{(\mathcal{N} \circ I)(\rho^{A^n B^n})} - \dim(B') \frac{\log n}{n} \right)$$

is achievable for supersymbols of length n.

$\Rightarrow R < I(B; B')_{(\mathcal{N} \circ I)(\rho^{A^n B^n})}$  is achievable.  $\square$

Converse:

Use reliable communication to construct correlated states  $(M, \hat{M})$  such that

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Use reliable communication to construct correlated states  $(M, \hat{M})$  such that

$$W^{M\hat{M}} \approx \sum_{i=1}^{2^{nR}} |i\rangle\langle i|^{\otimes n} \otimes |i\rangle\langle i|^{\hat{M}}$$
$$\triangleq \overline{\Phi}^{M\hat{M}}$$

$$\begin{aligned} \Rightarrow nR &= I(M; \hat{M})_{\overline{\Phi}} \\ &\leq I(M; \hat{M})_W + n\varepsilon' && \text{Fannes} \\ &\leq I(M; B^n \hat{B}^n) + n\varepsilon' && \text{DPI} \\ &\leq I(M B^n; B^n) + I(M; \hat{B}^n) + n\varepsilon' \\ &= I(M \hat{B}^n; B^n) + n\varepsilon' \\ &\leq I(N^{\otimes n}) + n\varepsilon' \quad \square \end{aligned}$$

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Stronger conclusion:

$$N + H(A)_{p^*} [qg] \geq I(N) [c \rightarrow c]$$

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