

Solution to Homework #7 (Seongwon Kang)

momentum eq.

$$\rho \left(u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right) = \rho_e U_e \frac{dU_e}{dx} + \frac{\partial}{\partial y} \left(\mu \frac{\partial u}{\partial y} \right)$$

$$\frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} = 0$$

B, C) $u(x, 0) = v(x, 0) = 0$

$$u(x, \infty) = U_e = U_0 \left(1 + x/4L \right)$$

I, E) $u(0, y) = U_{\text{uniform}}$

$$v(0, y) = 0$$

Let's nondimensionalize using following variables

$$\rho^* = \frac{\rho}{\rho_0}, \quad x^* = \frac{x}{L}, \quad y^* = \frac{y}{\delta}, \quad u^* = \frac{u}{U_0}$$

$$v^* = \frac{v}{U_0} \frac{L}{\delta}, \quad M^* = \frac{\mu}{\mu_0} = \frac{\mu}{\rho_0 U_0 \delta^2}$$

Then

$$\left(\begin{aligned} \rho^* \left(u^* \frac{\partial u^*}{\partial x^*} + v^* \frac{\partial u^*}{\partial y^*} \right) &= \rho_e^* U_e^* \frac{dU_e^*}{dx^*} + \frac{\partial}{\partial y^*} \left(M^* \frac{\partial u^*}{\partial y^*} \right) \\ \frac{\partial(\rho^* u^*)}{\partial x^*} + \frac{\partial(\rho^* v^*)}{\partial y^*} &= 0 \end{aligned} \right)$$

Using Busemann energy integral

$$h_0 = h + \frac{U^2}{2} = \text{const.}$$

$$\Rightarrow \underline{h^* + \frac{U^{*2}}{2} = \text{const.}}, \text{ where } h^* = \frac{h}{U_0^2}, U^* = \frac{U}{U_0}$$

From } h = C_p T

$$\Rightarrow \underline{h^* = \frac{C_p T_0}{U_0^2} T^*} \quad C_p = \frac{\gamma R}{\gamma - 1}$$

$$= \frac{\gamma R T_0}{U_0^2} \frac{1}{\gamma - 1} T^*$$

$$= \underline{\frac{1}{Ma^2} \frac{1}{\gamma - 1} T^*}$$

From } p = \rho R T

assuming $p(\mathcal{A}) \approx p_e$ since BL is thin,

$$\frac{p}{p_e} = \frac{T_e}{T} \Rightarrow \underline{\frac{p^*}{p_e^*} = \frac{T_e^*}{T^*}}$$

From } constant entropy s assumption

$$T ds = dh - v dp$$

$$\Rightarrow C_p dT = \frac{1}{\rho} dp$$

$$= R dT + \frac{R}{\rho} T dp$$

$$\Rightarrow C_v \frac{dT}{T} = R \frac{dp}{p} \Rightarrow T = C p^{\frac{R}{C_v}}$$

Note that the constant entropy assumption is only valid in the free-stream. In the boundary layer, pressure and density are obtained from the free-stream quantities using the ideal gas law and the fact that $p = p_e$.

$$\therefore \frac{T_1}{T_2} = \left(\frac{P_1}{P_2}\right)^{\kappa/\gamma} = \left(\frac{P_1}{P_2}\right)^{\gamma-1}$$

$$\Rightarrow \underline{T^* = p^{*\gamma-1}}$$

From $M = \beta T^n$

$$\frac{M}{M_0} = \frac{\beta T^n}{\beta T_0^n} \Rightarrow M^* = T^{*n}$$

where $M_0 = \beta T_0^n = \frac{\rho_0 U_0 \delta^2}{L}$

In summary,

$$\text{Compute } \left\{ \begin{array}{l} h^* \text{ from } h^* + \frac{U^{*2}}{2} = h_0^* \\ T^* \text{ from } h^* = \frac{1}{m_a^2} \frac{1}{\gamma-1} T^* \\ p^* \text{ from } p^* = T^{*\frac{1}{\gamma-1}} \text{ (for } p^* e^{\chi} \text{ ok.)} \\ \cdot \text{ from } \frac{p^*}{\rho e^*} = \frac{T e^*}{T^*} \\ M^* \text{ from } M^* = T^{*n} \end{array} \right.$$

Using implicit Euler method, from White

$$-\alpha^* U_{m+1,n+1} + (1 + 2\alpha^*) U_{m+1,n} - \alpha^* U_{m+1,n-1}$$

$$= U_{m,n} + \frac{\rho_{e,m} (U_{e,m+1}^2 - U_{e,m}^2)}{2\rho_{m,n} U_{m,n}} - (\beta - \gamma) (U_{m,n+1} - U_{m,n-1})$$

where

$$\alpha^* = \frac{M_{m,n} \Delta x}{\rho_{m,n} U_{m,n} \Delta y^2}, \quad \beta = \frac{\nu_{m,n} \Delta x}{2U_{m,n} \Delta y}, \quad \gamma = \frac{(M_{m,n+1} - M_{m,n-1}) \Delta x}{4\rho_{m,n} U_{m,n} \Delta y^2}$$

Computational procedure

- 1) From U_e , compute T_e , P_e , M_e
- 2) Marching in x using discretized momentum eq. to get U at $m+1$
- 3) compute T_{m+1} , P_{m+1} , M_{m+1}
- 4) compute V_{m+1} using continuity eq.

Test case: Blasius case $U_e = U_0 = \text{const.}$

$$Ma = 1.5, \quad \gamma = 1.4 \quad \text{for air, } n = 1$$

for nondimensionalization, following relation is used

$$x^* = \frac{x}{L}, \quad y^* = \frac{y}{\delta}, \quad u^* = \frac{u}{U_0}, \quad v^* = \frac{v}{U_0} \frac{L}{\delta}$$

Domain size in x : $L_x = \underline{2\delta}$

in y : $L_y = \underline{4\delta}$

grid points in x : $N_x = \underline{800}$

in y : $N_y = \underline{800}$

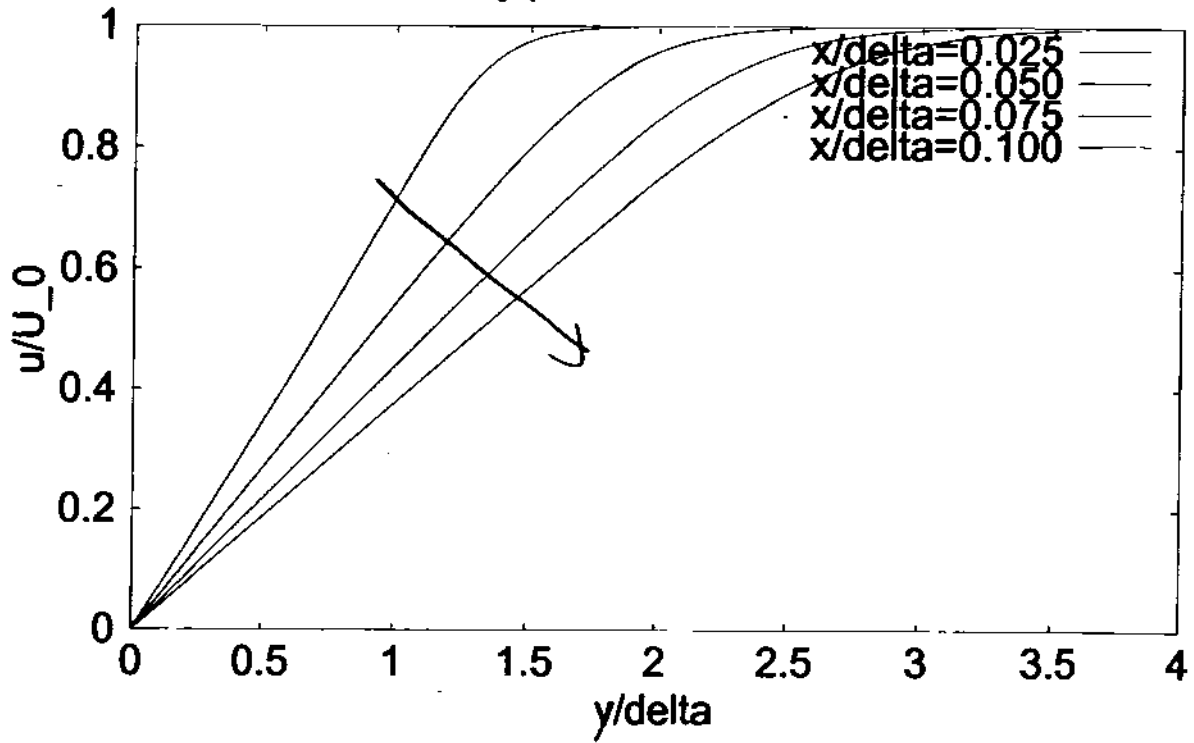
Corrected
mistake in
Seongwon's
solution

(3/10/04):

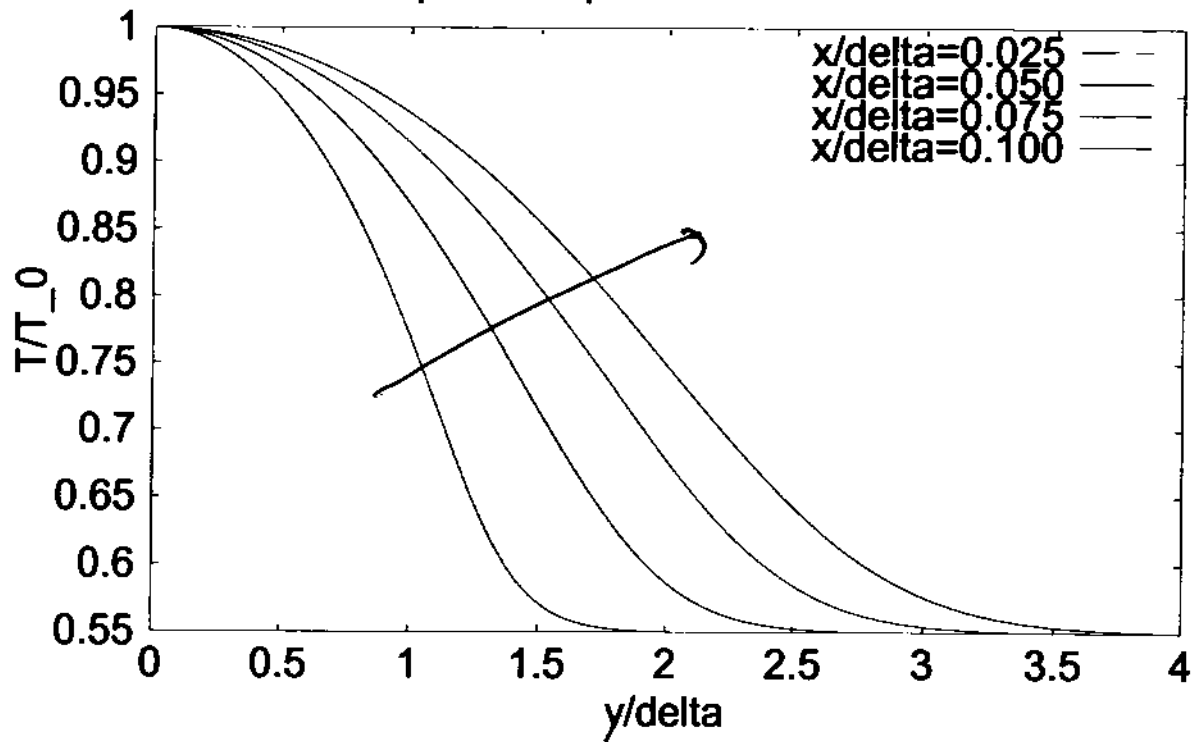
$$x^* = x/L$$

$$v^* = (v/U_0)(L/\delta)$$

Velocity profiles as downstream



Temperature profiles as downstream



Profiles are developing as downstream.

Main case

$$U_e = U_0 \left(1 + \frac{x}{4L} \right) \Rightarrow U_e^* = 1 + \frac{x^*}{4}$$

$$Ma = 1.5, \quad \gamma = 1.4, \quad n = 2/3$$

$$\text{Domain size in } x : L_x = \underline{4L}$$

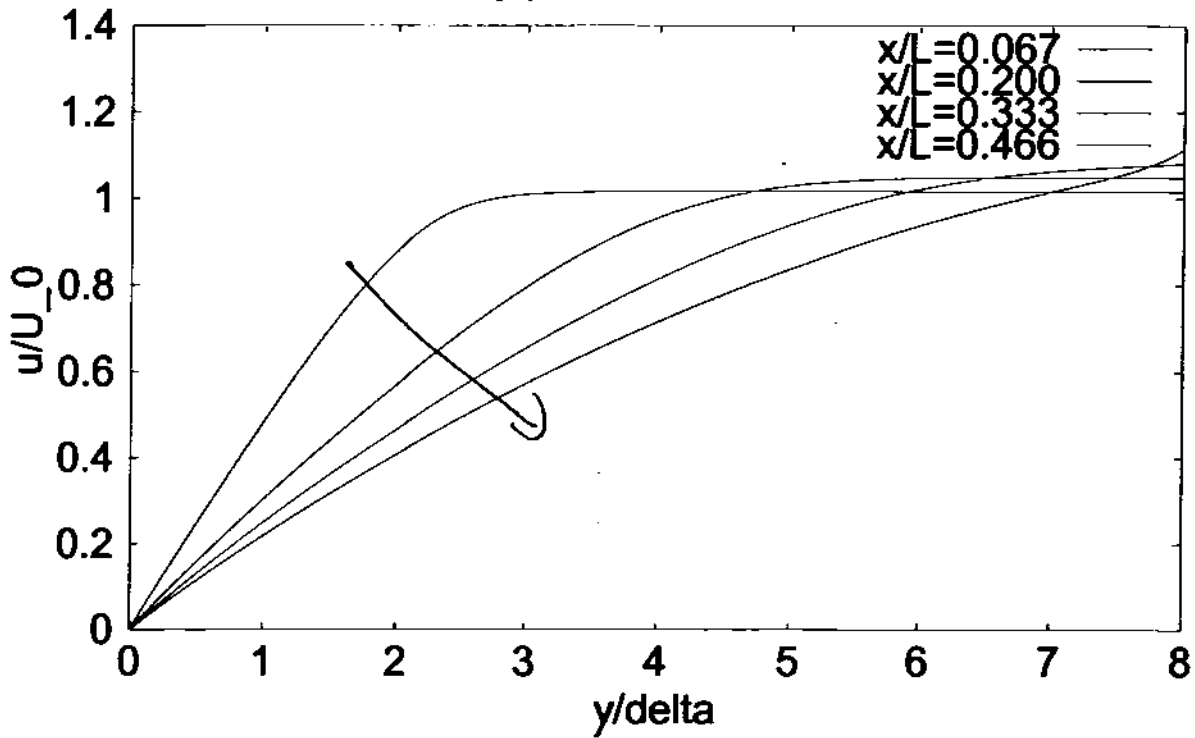
$$\text{in } y : L_y = \underline{8\delta}$$

$$\# \text{ grid points in } x : N_x = \underline{1200}$$

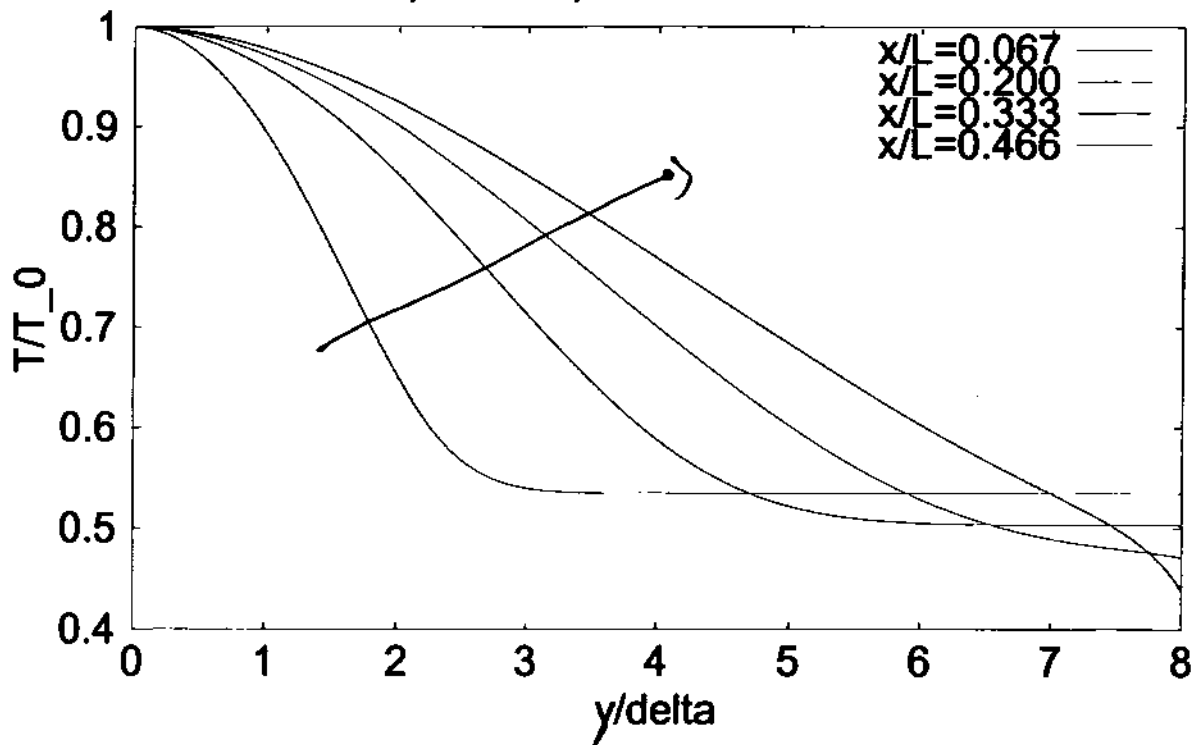
$$N_y = \underline{1200}$$

In this case, U_e is developing as downstream and at certain x^* , T^* will become negative from $h^* + \frac{U^{*2}}{2} = \text{const.}$ then further marching will be impossible

Velocity profiles as downstream



Temperature profiles as downstream



Unlike Blasius case, farfield velocity increases and farfield temperature decreases as downstream.

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C-----
C FORTRAN CODE FOR COMPRESSIBLE BOUNDARY LAYER FLOW
C FOR ME 351B
C
C-----
C MADE BY SEONGWON KANG
C-----
PROGRAM MAIN
PARAMETER(M_X=1500)
PARAMETER(M_Y=1300)
REAL X(0:M_X+1),Y(0:M_Y+1)
REAL U(0:M_X+1,0:M_Y+1),V(0:M_X+1,0:M_Y+1)
REAL T(0:M_X+1,0:M_Y+1),RHO(0:M_X+1,0:M_Y+1),VSC(0:M_X+1,0:M_Y+1)
REAL, DIMENSION (M_Y) :: AU,BU,CU,UU
INTEGER IP0S(6)
CHARACTER*20 FNAME(6)
! INPUT PARAMETERS
! -----
X_L=4.
Y_L=8.
N_X=1199
N_Y=1199
! -----
! INITIALIZATION OF GRID
! -----
DX=X_L/(N_X+1)
DY=Y_L/(N_Y+1)
DO I=0,N_X+1
  X(I)=I*DX
END DO
DO J=0,N_Y+1
  Y(J)=J*DY
END DO
JE=N_Y+1
! -----
! INITIAL & BOUNDARY VELOCITY
! -----
U=0.; V=0.
C DO I=0,N_X+1
C U(I,JE)=1.
C END DO
DO I=0,N_X+1
  U(I,JE)=1+X(I)/4
END DO
DO J=0,N_Y
  U(0,J)=U(0,JE)

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V(0,J)=0.
END DO
! -----
! COMPRESSIBLE PROPERTIES
! -----
GAMMA=1.4
V_MACH=1.5
C POWER=1.
C BLASIUS
POWER=2./3.
ACCELERATING FLOW
T_INF=1.
H_STAG=T_INF/(V_MACH**2*(GAMMA-1))
! -----
! INITIAL & BOUNDARY COMPRESSIBLE PROPERTIES
! -----
DO I=0,N_X+1
  T(I,JE)=(H_STAG-0.5*U(I,JE)**2)*V_MACH**2*(GAMMA-1)
  RHO(I,JE)=T(I,JE)**(1./(GAMMA-1))
END DO
DO J=0,N_Y+1
  T(0,J)=(H_STAG-0.5*U(0,J)**2)*V_MACH**2*(GAMMA-1)
  RHO(0,J)=RHO(0,JE)*T(0,JE)/T(0,J)
  VSC(0,J)=T(0,J)**POWER
END DO
! -----
! MARCHING IN X
! -----
DO I=1,N_X
  IM=I-1
  ! COMPUTATION OF U
  DO J=1,N_Y
    ALP=VSC(IM,J)*DX/(RHO(IM,J)*U(IM,J)*DY**2)
    BET=V(IM,J)*DX/(2*U(IM,J)*DY)
    SIG=(VSC(IM,J+1)-VSC(IM,J-1))*DX/
      (4*RHO(IM,J)*U(IM,J)*DY**2)
    AU(J)=-ALP
    BU(J)=-ALP
    BU(J)=1+2*ALP
    RU(J)=U(IM,J)+RHO(IM,JE)*(U(I,JE)**2-U(IM,JE)**2)/
      (2*RHO(IM,J)*U(IM,J))-
      (BET-SIG)*(U(IM,J+1)-U(IM,J-1))
  END DO
  RU(1)=RU(1)-AU(1)*U(I,0)
  RU(N_Y)=RU(N_Y)-CU(N_Y)*U(I,JE)
CALL TRDIAC_A(M_Y,1,AU,BU,CU,RU,UU,1,N_Y,1,1)

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DO J=1,N,Y
U(I,J)=UU(J)
END DO

! COMPUTATION OF T,RHO,VSC
DO J=0,N,Y+1
T(I,J)=(H_STAG-0.5*U(I,J)**2)*V MACH**2*(GAMMA-1)
RHO(I,J)=RHO(I,JE)*T(I,JE)/T(I,J)
VSC(I,J)=T(I,J)**POWER
END DO

! COMPUTATION OF V
DO J=1,N,Y+1
V(I,J)=RHO(I,J-1)*V(I,J-1)/RHO(I,J)-DY/(2*RHO(I,J)*DX)*
( RHO(I,J)*U(I,J)-RHO(IM,J)*U(IM,J)+
RHO(I,J-1)*U(I,J-1)-RHO(IM,J-1)*U(IM,J-1) )
END DO

END DO
! -----
! OUTPUT
IPOS(1)=20 ; FNAME(1)='V1'
IPOS(2)=60 ; FNAME(2)='V2'
IPOS(3)=100 ; FNAME(3)='V3'
IPOS(4)=140 ; FNAME(4)='V4'

DO II=1,4
OPEN(3,FILE=FNAME(II))
DO J=0,N,Y+1
WRITE(3,88)Y(J),U(IPOS(II),J),V(IPOS(II),J),T(IPOS(II),J)
END DO
CLOSE(3)
END DO

88 FORMAT(4E16.8)

STOP
END

C***** TRDIAG A *****
C SOLVE TRIDIAGONAL MATRICES WITH
C DIFFERENT COEFS.

SUBROUTINE TRDIAG_A(M1,M2,A,B,C,R,UU,L1,L2,LL1,LL2)
REAL, DIMENSION(M2,M1) :: A,B,C,R,UU,GAM
REAL, DIMENSION(M2) :: BET
DO I=LL1,LL2
BET(I)=B(I,L1)

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UU(I,LL1)=R(I,LL1)/BET(I)
END DO

DO J=LL1+1,LL2
DO I=LL1,LL2
GAM(I,J)=C(I,J-1)/BET(I)
BET(I)=B(I,J)-A(I,J)*GAM(I,J)
UU(I,J)=(R(I,J)-A(I,J)*UU(I,J-1))/BET(I)
END DO
END DO

DO J=LL2-1,LL1,-1
DO I=LL1,LL2
UU(I,J)=UU(I,J)-GAM(I,J+1)*UU(I,J+1)
END DO
END DO

RETURN
END

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