Lecture 7 ELE 301: Signals and Systems

Prof. Paul Cuff

Princeton University

Fall 2011-12

Introduction to Fourier Transforms

- · Fourier transform as a limit of the Fourier series
- Inverse Fourier transform: The Fourier integral theorem
 - · Example: the rect and sinc functions
 - Cosine and Sine Transforms
 - · Symmetry properties
 - ullet Periodic signals and δ functions

Cuff (Lecture 7) ELE 201: Signals and Systems Fall 2011-12 1/22

101100121121 2 000

Fourier Series

Suppose x(t) is not periodic. We can compute the Fourier series as if x was periodic with period T by using the values of x(t) on the interval $t \in [-T/2, T/2)$.

$$a_k = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-j2\pi k l_0 t} dt,$$

 $x_T(t) = \sum_{k=-\infty}^{\infty} a_k e^{j2\pi k l_0 t},$

where $f_0 = 1/T$.

The two signals x and x_T will match on the interval [-T/2, T/2) but $\bar{x}(t)$ will be periodic.

What happens if we let T increase?

Fall 2011-12 3 / 22

Rect Example

For example, assume $x(t)=\mathrm{rect}(t)$, and that we are computing the Fourier series over an interval T,



The fundamental period for the Fourier series in $\,T_{\rm o}$, and the fundamental frequency is $\,f_0=1/T_{\rm o}$.

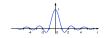
The Fourier series coefficients are

$$a_k = \frac{1}{\tau} \operatorname{sinc}(kf_0)$$

where $sinc(t) = \frac{sin(\pi t)}{\pi t}$.

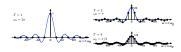


The Sinc Function



Rect Example Continued

Take a look at the Fourier series coefficients of the rect function (previous slide). We find them by simply evaluating $\frac{1}{T}\mathrm{sinc}(f)$ at the points $f=kf_0$.



More densely sampled, same sinc() envelope, decreased amplitude.

off (Lecture 7) ELECOL: Signals and Systems Fall 2011-12 6/22

Ceff (Lecture 7) ELE 301: Signals and Systems Fall 2011-12 5 / 22

Fourier Transforms

Given a continuous time signal x(t), define its Fourier transform as the function of a real f:

$$X(f) = \int_{-\infty}^{\infty} x(t)e^{-j2\pi ft} dt$$

This is similar to the expression for the Fourier series coefficients.

Note: Usually X(f) is written as $X(i2\pi f)$ or $X(i\omega)$. This corresponds to the Laplace transform notation which we encountered when discussing transfer functions H(s).

We can interpret this as the result of expanding x(t) as a Fourier series in an interval [-T/2,T/2), and then letting $T\to\infty$.

The Fourier series for x(t) in the interval [-T/2, T/2):

$$x_T(t) = \sum_{k=-\infty}^{\infty} a_k e^{j2\pi k f_0 t}$$

where

$$a_k = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-j2\pi k f_0 t} dt.$$

Define the truncated Fourier transform:

$$X_T(f) = \int_{-\frac{T}{2}}^{\frac{T}{2}} x(t)e^{-j2\pi ft} dt$$

so that

$$a_k = \frac{1}{T} X_T(kf_0) = \frac{1}{T} X_T\left(\frac{k}{T}\right).$$

The Fourier series is then

$$x_T(t) = \sum_{k=-\infty}^{\infty} \frac{1}{T} X_T(kf_0) e^{j2\pi kf_0 t}$$

The limit of the truncated Fourier transform is

$$X(f) = \lim_{T \to \infty} X_T(f)$$

The Fourier series converges to a Riemann integral:

$$\begin{array}{ll} x(t) & = & \lim_{T \to \infty} x_T(t) \\ & = & \lim_{T \to \infty} \sum_{k = -\infty}^{\infty} \frac{1}{T} X_T\left(\frac{k}{T}\right) e^{j2\pi \frac{k}{T}t} \\ & = & \int_{-\infty}^{\infty} X(f) e^{j2\pi ft} df. \end{array}$$

Continuous-time Fourier Transform

Which yields the inversion formula for the Fourier transform, the Fourier integral theorem:

$$X(f) = \int_{-\infty}^{\infty} x(t)e^{-j2\pi ft} dt,$$

$$x(t) = \int_{-\infty}^{\infty} X(f)e^{j2\pi ft} df.$$

Cell (Lecture 2) ELECTRIC Strain and Support

Comments.

- There are usually technical conditions which must be satisfied for the integrals to converge – forms of smoothness or Dirichlet conditions.
- The intuition is that Fourier transforms can be viewed as a limit of Fourier series as the period grows to infinity, and the sum becomes an integral.
- $\int_{-\infty}^{\infty} X(f)e^{j2\pi ft} df$ is called the *inverse Fourier transform* of X(f). Notice that it is identical to the Fourier transform except for the sign in the exponent of the complex exponential.
- If the inverse Fourier transform is integrated with respect to ω rather than f, then a scaling factor of 1/(2π) is needed.

Cosine and Sine Transforms

Assume x(t) is a possibly complex signal.

$$\begin{split} X(f) &= \int_{-\infty}^{\infty} x(t) e^{-j2\pi it} dt \\ &= \int_{-\infty}^{\infty} x(t) \left(\cos(2\pi f t) - j \sin(2\pi f t) \right) dt \\ &= \int_{-\infty}^{\infty} x(t) \cos(\omega t) dt - j \int_{-\infty}^{\infty} x(t) \sin(\omega t) dt. \end{split}$$

Fourier Transform Notation

For convenience, we will write the Fourier transform of a signal x(t) as

$$\mathcal{F}[x(t)] = X(f)$$

and the inverse Fourier transform of X(f) as

$$\mathcal{F}^{-1}[X(f)] = x(t).$$

Note that

$$\mathcal{F}^{-1}\left[\mathcal{F}\left[x(t)\right]\right]=x(t)$$

and at points of continuity of x(t).

Duality

Notice that the Fourier transform \mathcal{F} and the inverse Fourier transform \mathcal{F}^{-1} are almost the same

Duality Theorem: If $x(t) \Leftrightarrow X(f)$, then $X(t) \Leftrightarrow x(-f)$.

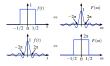
In other words, $\mathcal{F}[\mathcal{F}[x(t)]] = x(-t)$.

Example of Duality

• Since $rect(t) \Leftrightarrow sinc(f)$ then

$$\operatorname{sinc}(t) \Leftrightarrow \operatorname{rect}(-f) = \operatorname{rect}(f)$$

(Notice that if the function is even then duality is very simple)



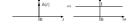
A unit impulse $\delta(t)$ is not a signal in the usual sense (it is a generalized function or distribution). However, if we proceed using the sifting property, we get a result that makes sense:

Generalized Fourier Transforms: & Functions

$$\mathcal{F}\left[\delta(t)\right] = \int_{-\infty}^{\infty} \delta(t)e^{-j2\pi ft} dt = 1$$

$$\delta(t) \Leftrightarrow 1$$

This is a generalized Fourier transform. It behaves in most ways like an ordinary FT.



Shifted δ

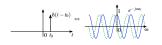
A shifted delta has the Fourier transform

$$\mathcal{F}[\delta(t-t_0)] = \int_{-\infty}^{\infty} \delta(t-t_0)e^{-j2\pi ft} dt$$

$$= e^{-j2\pi t_0 f}$$

so we have the transform pair

$$\delta(t - t_0) \Leftrightarrow e^{-j2\pi t_0 f}$$



Constant

Next we would like to find the Fourier transform of a constant signal x(t) = 1. However, direct evaluation doesn't work:

$$\mathcal{F}[1] = \int_{-\infty}^{\infty} e^{-j2\pi ft} dt$$

= $\frac{e^{-j2\pi ft}}{-j2\pi f}\Big|_{-\infty}^{\infty}$

and this doesn't converge to any obvious value for a particular f.

We instead use duality to guess that the answer is a δ function, which we can easily verify.

$$\mathcal{F}^{-1}\left[\delta(f)\right] = \int_{-\infty}^{\infty} \delta(f)e^{j2\pi ft} df$$
$$= 1.$$

So we have the transform pair



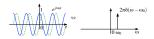
This also does what we expect – a constant signal in time corresponds to an impulse a zero frequency.

Sinusoidal Signals

If the δ function is shifted in frequency,

$$\begin{split} \mathcal{F}^{-1}\left[\delta(f-f_0)\right] &= \int_{-\infty}^{\infty} \delta(f-f_0) e^{j2\pi f_0} df \\ &= e^{j2\pi f_0 t} \end{split}$$

 $e^{j2\pi f_0 t} \Leftrightarrow \delta(f - f_0)$



Cosine

With Euler's relations we can find the Fourier transforms of sines and cosines

$$\begin{split} \mathcal{F}[\cos(2\pi f_0 t)] &= \mathcal{F}\left[\frac{1}{2}\left(e^{j2\pi f_0 t} + e^{-j2\pi f_0 t}\right)\right] \\ &= \frac{1}{2}\left(\mathcal{F}\left[e^{j2\pi f_0 t}\right] + \mathcal{F}\left[e^{-j2\pi f_0 t}\right]\right) \\ &= \frac{1}{2}\left(\delta(f - f_0) + \delta(f + f_0)\right). \end{split}$$

30

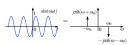
$$cos(2\pi f_0 t) \Leftrightarrow \frac{1}{2} (\delta(f - f_0) + \delta(f + f_0)).$$



Sine

Similarly, since $sin(f_0t) = \frac{1}{77}(e^{j2\pi f_0t} - e^{-j2\pi f_0t})$ we can show that

$$sin(f_0t) \Leftrightarrow \frac{j}{2}(\delta(f + f_0) - \delta(f - f_0)).$$



The Fourier transform of a sine or cosine at a frequency f_0 only has energy exactly at $\pm f_0$, which is what we would expect.

EN 2011-12 22 / 22